EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH (CERN)

CERN-EP-2022-001
2022/03/01

CMS-B2G-20-002

# Search for a W′ boson decaying to a vector-like quark and a top or bottom quark in the all-jets final state at $\sqrt{s} = 13\,\text{TeV}$

The CMS Collaboration*


## Abstract

A search is presented for a heavy W′ boson resonance decaying to a B or T vector-like quark and a t or a b quark, respectively. The analysis is performed using proton-proton collisions collected with the CMS detector at the LHC. The data correspond to an integrated luminosity of $138\,\text{fb}^{-1}$ at a center-of-mass energy of 13 TeV. Both decay channels result in a signature with a t quark, a Higgs or Z boson, and a b quark, each produced with a significant Lorentz boost. The all-hadronic decays of the Higgs or Z boson and of the t quark are selected using jet substructure techniques to reduce standard model backgrounds, resulting in a distinct three-jet W′ boson decay signature. No significant deviation in data with respect to the standard model background prediction is observed. Upper limits are set at 95% confidence level on the product of the W′ boson cross section and the final state branching fraction. A W′ boson with a mass below 3.1 TeV is excluded, given the benchmark model assumption of democratic branching fractions. In addition, limits are set based on generalizations of these assumptions. These are the most sensitive limits to date for this final state.


*Submitted to the Journal of High Energy Physics*



*See Appendix A for the list of collaboration members

# 1 Introduction

New massive charged gauge bosons are predicted by various extensions of the standard model (SM) [1–3]. The W′ boson is a hypothetical heavy partner of the SM W gauge boson that could be produced in proton-proton (pp) collisions at the CERN LHC. Searches for W′ bosons have most recently been performed at a center-of-mass energy of 13 TeV by the CMS and ATLAS Collaborations in the lepton-neutrino [4, 5], diboson [6, 7], and diquark [8, 9] final states. The decay of the W′ boson to a heavy B or T vector-like quark (VLQ) and a t or b quark, respectively, is predicted, e.g., in composite Higgs boson models with custodial symmetry protection [10–12]. The VLQs are hypothetical heavy partners of SM quarks that contain left- and right-handed chiralities, which transform identically under the SM gauge groups. These models stabilize the quantum corrections to the Higgs boson mass and preserve naturalness. The W′ branching fraction to a quark and a VLQ depends on the VLQ mass, with a maximum of 50% [13]. Searches for VLQs have been performed by CMS and ATLAS in both the single- [14–17] and pair-production [18–21] channels.

This paper describes a search for the W′ boson in the all-hadronic final state, where the B VLQ decays into a Higgs or Z boson and a b quark, or the T VLQ decays into a Higgs or Z boson and a t quark. The separate studies of the two decay channels are referred to as the tHb analysis and the tZb analysis throughout the paper. Both the B- and T-mediated decays result in the same signature, as shown in Fig. 1.

Owing to the high W′ boson and VLQ masses considered, the decay products are highly Lorentz boosted, and are combined into large-radius jets with distinct substructure. The challenge is to distinguish between these events and the main background processes, which arise from events comprised of light quark and gluon jets produced via the strong interaction, referred to as quantum chromodynamics (QCD) multijet events, and t quark pair ($\mathrm{t\bar{t}}$) events. These backgrounds are modeled using a combination of Monte Carlo simulation and control regions in data. The invariant mass distribution of the three-jet system ($m_{\mathrm{tHb}}$ or $m_{\mathrm{tZb}}$) is used to search for the W′ boson signal. The data correspond to an integrated luminosity of 138 $\mathrm{fb}^{-1}$ [22–24] of pp collisions at $\sqrt{s} = 13$ TeV, recorded in the years 2016 to 2018.

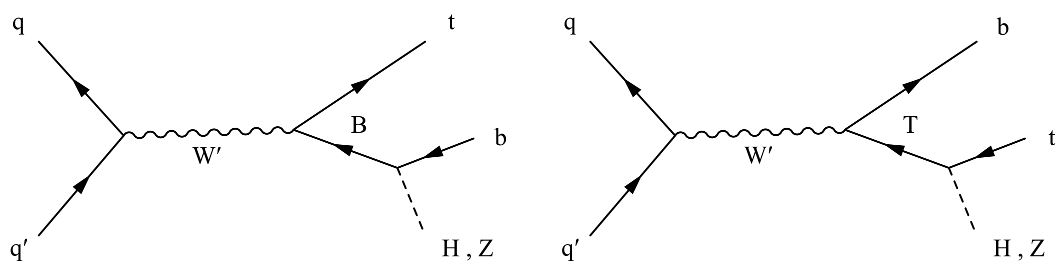


Figure 1: Dominant Feynman diagrams for the signal model considered.

We consider a theoretical framework where the t quark and W′ boson are superpositions of elementary and composite modes, with the t quark degree of compositeness given by $s_{\mathrm{L}}$, and the mixing angle of the elementary and composite W′ states given by $\theta_2$ [13]. The W′ boson is assumed to be produced in a Drell–Yan process, with a cross section that is inversely proportional to $\cot^2(\theta_2)$. Small values of $\cot(\theta_2)$ tend to be dominated by the leptonic W′ boson decay mode. Large values of the $s_{\mathrm{L}}$ parameter increase the relative phase space for the decay into two VLQs, whereas small $s_{\mathrm{L}}$ values enhance the W′ diboson decays. The analysis assumes this theoretical framework as evaluated at $s_{\mathrm{L}} = 0.5$ and $\cot(\theta_2) = 3$, which is chosen to enhance sensitivity to the single VLQ decay channel of the W′ boson. The expected signal cross sections are evaluated at $\sqrt{s} = 13$ TeV, considering W′ boson masses in the 1.5 to 5.0 TeV range with the assumption that the W′ → VLQ branching fraction is equally distributed between the tB and tT final states. For the benchmark, the VLQ branching fractions for the decays to qH and qZ are both assumed to be 50%.

Tabulated results are provided in the HEPData record for this analysis [25].

## 2 The CMS detector

The central feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter (ECAL), and a brass and scintillator hadron calorimeter (HCAL), each composed of a barrel and two endcap sections. Forward calorimeters extend the pseudorapidity ($\eta$) coverage provided by the barrel and endcap detectors. Muons are detected in gas-ionization chambers embedded in the steel flux-return yoke outside the solenoid.

Events of interest are selected using a two-tiered trigger system. The first level, composed of custom hardware processors, uses information from the calorimeters and muon detectors to select events at a rate of around 100 kHz within a fixed latency of about 4 $\mu$s [26]. The second level, known as the high-level trigger, consists of a farm of processors running a version of the full event reconstruction software optimized for fast processing, and reduces the event rate to around 1 kHz before data storage [27].

A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [28].

The particle-flow algorithm [29] aims to reconstruct and identify each individual particle with an optimized combination of information from the various elements of the CMS detector. The energy of each photon is obtained from the ECAL measurement. The energy of each electron is determined from a combination of the electron momentum at the primary interaction vertex as determined by the tracker, the energy of the corresponding ECAL cluster, and the energy sum of all bremsstrahlung photons spatially compatible with originating from the electron track. The energy of each muon is obtained from the momentum, which is measured by the curvature of the corresponding track. The energy of charged hadrons is determined from a combination of their momentum measured in the tracker and the matching ECAL and HCAL energy deposits, corrected for zero-suppression effects and for the response function of the calorimeters to hadronic showers. Finally, the energy of each neutral hadron is obtained from the corresponding corrected ECAL and HCAL energies that are not associated with a charged hadron track.

Jets are clustered with the anti-$k_{\mathrm{T}}$ algorithm [30] using the FASTJET 3.0 [31] software package with distance parameters of 0.8 (AK8 jets) and 0.4 (AK4 jets). Jet momentum is determined as the vectorial sum of all particle momenta in the jet, and is found from simulation to be within 5 to 10% of the true momentum over the entire $p_{\mathrm{T}}$ spectrum and detector acceptance. Additional pp interactions within the same or nearby bunch crossings (pileup) can contribute additional tracks and calorimetric energy depositions to the jet momentum. To mitigate this effect, tracks identified to be originating from pileup vertices are discarded and an offset correction is applied to correct for remaining contributions.

The candidate vertex with the largest value of summed physics-object $p_{\mathrm{T}}^2$ is taken to be the primary pp interaction vertex. The physics objects used for this determination are the jets, clustered using the jet finding algorithm [30, 31] with the tracks assigned to candidate vertices as inputs, and the associated missing transverse momentum, taken as the negative vector sum of the $p_T$ of those jets.

To account for the neutral pileup component of the AK8 jets, the pileup per particle identification (PUPPI) algorithm [32] is used, which applies weights that rescale the jet $p_T$ based on the per-particle probability of originating from the primary vertex prior to jet clustering.

Jet energy corrections are derived from simulation studies so that the average measured response of jets becomes identical to that of particle level jets. In situ measurements of the momentum balance in dijet, photon + jet, Z + jet, and multijet events are used to determine any residual differences between the jet energy scale (JES) in data and in simulation, and appropriate corrections are made [33]. Additional selection criteria are applied to each jet to remove jets potentially dominated by instrumental effects or reconstruction failures. The jet energy resolution (JER) amounts typically to 15% at 10 GeV, 8% at 100 GeV, and 4% at 1 TeV, to be compared to about 40, 12, and 5% obtained when the calorimeters alone are used for jet clustering.

## 3 Simulated samples

The $t\bar{t}$ and single t quark production backgrounds are estimated from simulation generated at next-to-leading order with POWHEG 2.0 [34–39]. The signal and QCD background events are generated at leading order using the MADGRAPH5_aMC@NLO 2.3.3 [40, 41] package. Signal events are generated in the mass range from 1.5 to 5.0 TeV in 0.5 TeV increments. Events are generated using either the NNPDF 3.0 [42] or 3.1 [43] parton distribution functions (PDFs), with parton showering and hadronization simulated with PYTHIA 8.212 [44].

The simulated events are produced with either the CUETP8M2T4 [45], CUETP8M1 [46], or CP5 [47] underlying-event tunes. The CMS detector is simulated using GEANT4 [48]. All simulated events include pileup simulation and are weighted such that the distribution of the number of interactions per bunch crossing agrees with that in data assuming a total inelastic cross section of 69.2 mb.

For each W′ boson mass ($m_{W'}$) point, three characteristic VLQ mass points are generated, corresponding to low (1/2 $m_{W'}$), medium (2/3 $m_{W'}$), and high (3/4 $m_{W'}$) mass ratios, in order to allow the full sensitive phase space of the boosted W′ boson decay products to be explored.

The generated W′ boson and VLQ widths are chosen to be narrow (less than 10%) relative to the experiment resolution, as required by the theoretical predictions for the region of phase space under investigation [13].

## 4 Event reconstruction

The final state is characterized by three high-$p_T$ jets from the decay of the t and b quarks, and from either the Higgs or Z boson. Because the signal of interest corresponds to a high mass resonance decaying to multiple high-$p_T$ jets, data events are triggered by the presence of either a large value of $H_T$ (the scalar sum of all AK4 jet $p_T$ in the event), an AK8 jet with a large $p_T$, or an AK8 jet with a large jet mass.

The efficiency of the trigger selection is studied using a sample of events that are triggered by at least one high-$p_T$ muon. The fraction of these events that pass the full trigger selection is defined as the trigger efficiency. The offline event selection requires that the $H_T$ is larger than 1 TeV and the maximum jet mass is larger than 140 GeV, which ensures that the trigger is close to fully efficient. The small inefficiency (less than 1%) due to the trigger selection is taken into account as an event weight when processing simulated events.

The jets from the t quark (t jet) and from the Higgs or Z boson (Higgs or Z jet) decay tend to be wide, massive, and have a distinct substructure, whereas the jet from the b quark (b jet) tends to be narrower and have a lower mass. Therefore, an AK8 jet with $p_{\mathrm{T}} > 400\,\mathrm{GeV}$ is required for the Higgs or Z boson candidate, and similarly for the t quark candidate, together with an AK4 jet with $p_{\mathrm{T}} > 200\,\mathrm{GeV}$ for the b candidate. The angular separation $\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2}$, where $\phi$ is the azimuthal angle, between the two AK8 jets is required to be at least 1.6 to reduce the confusion of jet shapes arising from the abutment of jet boundaries, which can bias the background estimate. The AK8 jets are then selected as being consistent with a t quark or a Higgs or Z boson decay using the jet identification ("tagging") procedures defined below. The jet considered for the b quark candidate is the leading AK4 jet with a $\Delta R$ of at least 1.2 from the tagged AK8 jets. All jets are required to be within the tracker acceptance ($|\eta| < 2.4$).

## 4.1 Tagging t jets

For t quarks with $p_{\mathrm{T}} > 400\,\mathrm{GeV}$, the decay products, one b quark and two light quarks, can merge into a single AK8 jet. The t quark jets are identified by a novel algorithm designed to separate t jets from jets arising from light quarks and gluons using a deep convolutional neural network called imageTop [49]. The network uses the particle momentum and particle-flow identification category as an analogy to the pixel intensity and color that would be used as inputs to a conventional image-classification algorithm. Additionally, the algorithm includes b tagging information simultaneously in the training in order to identify the b quark from the t quark decay.

The version used in this analysis includes a network that is trained on QCD jets that are drawn from a mass distribution that is constrained to be identical to the top mass distribution, so as to decorrelate the discriminator from the jet mass (imageTop$_{\mathrm{MD}}$). This analysis is the first application of the imageTop$_{\mathrm{MD}}$ tagger, which is the most sensitive decorrelated t tagger developed by the CMS Collaboration, and the mass decorrelation is valid over a large range of t jet mass and $p_{\mathrm{T}}$. This tagger shows an improvement of approximately a factor of six in background rejection at a constant efficiency when compared to the best t quark taggers used previously. The imageTop$_{\mathrm{MD}}$ tagger is used to select t quark jets at an operating point with $\approx$60% efficiency and a mistag probability of $\approx$2% in QCD multijet events.

Owing to the decorrelated nature of the t tagging algorithm, we also apply a selection on the mass of the t jet. The modified mass-drop tagger algorithm [50], also known as the "soft drop" algorithm [51], with $\beta = 0$ and $z = 0.1$, is used to calculate the mass variable $m_{\mathrm{SD}}(\mathrm{t})$ to discriminate between t jets and the typically lighter background jets. This variable is corrected by applying AK4 jet energy corrections to the subjets identified by the algorithm. For events in the signal region, we require $140 < m_{\mathrm{SD}}(\mathrm{t}) < 220\,\mathrm{GeV}$, to be consistent with the reconstructed mass of a t quark.

A simulation-to-data scale factor for the t tagging requirement is extracted by fitting the soft-drop mass distribution in a highly enriched sample of SM $\mathrm{t\bar{t}}$ events in data and simulation. This scale factor is applied to correct differences between the top tagging efficiencies for simulated events and data. The scale factor is $p_{\mathrm{T}}$ dependent with a typical value of $\approx$0.95 and does not vary significantly between different running periods.

## 4.2 Tagging Higgs jets

In the case of a highly boosted Higgs boson in the $b\overline{b}$ decay mode, the decay products tend to merge into a single jet with a mass that is consistent with that of a Higgs boson and contains two b hadron decays. The soft-drop algorithm is used to extract the variable $m_{\mathrm{SD}}(\mathrm{H})$ as a measure of the jet mass, but in this case, the jet is scaled using a $p_{\mathrm{T}}$- and $\eta$-dependent correction suitable for two-prong resonances below the t quark mass [52].

To be selected, the jet should fulfill the condition $105 < m_{\mathrm{SD}}(\mathrm{H}) < 140\,\mathrm{GeV}$, to be consistent with the reconstructed Higgs boson mass. Scale factors are used for the jet mass scale and resolution (JMS and JMR), which are derived from a fit to the soft-drop mass distribution of the W boson in a sample enriched in lepton+jets $t\overline{t}$ events using the technique outlined in Ref. [53].

To identify the two b hadrons clustered into the merged Higgs jet, a dedicated double-b tagging algorithm (Dbtag) is used at an operating point that has an efficiency of ≈60% and a mistag probability of ≈7% for u, d, s quarks and gluon jets. This algorithm aims to identify two b quarks from the Higgs decay by using variables correlated with b hadron lifetime and mass combined using a boosted decision tree. Data enriched in QCD-produced $b\overline{b}$ and $t\overline{t}$ events are used to establish scale factors for Dbtag for signal and mistagged t quarks, respectively [54].

Figure 2 shows the simulated distributions of discriminants that are used for t and Higgs jet tagging in $t\overline{t}$, QCD, and signal for the tHb analysis.

## 4.3 Tagging Z jets

In contrast to the Higgs boson, the Z boson has significant decay branching fractions to all five kinematically accessible quark flavors, so we require a jet energy deposition pattern that is consistent with coming from two hard partons, along with a jet mass that is consistent with a Z boson.

The $N$-subjettiness [55] algorithm defines the $\tau_N$ variable, which quantifies the consistency of the jet energy pattern with the expectation from $N$ or fewer hard partons. The $N$-subjettiness is defined as

$$\tau_N = \frac{1}{d} \sum_i p_{\mathrm{T}i} \min\{\Delta R_{1,i}, \Delta R_{2,i}, \ldots, \Delta R_{N,i}\}, \tag{1}$$

where the sum is over jet constituents. The $\tau_N$ discriminant is bound between zero and one, with low values being more consistent with $N$ or fewer partons. For a Z boson hadronic decay, the ratio of $\tau_2$ to $\tau_1$ ($\tau_{21}$) is used at an operating point of ≈85% efficiency and a mistag probability of ≈20% in QCD multijet events. Scale factors used for the $N$-subjettiness selection are derived simultaneously with those for the JMS and JMR as discussed in Section 4.2.

The $m_{\mathrm{SD}}(\mathrm{Z})$ is required to be within the range $65 < m_{\mathrm{SD}}(\mathrm{Z}) < 105\,\mathrm{GeV}$, to be consistent with the mass of a Z boson.

Figure 3 shows the distributions of discriminants that are used for t and Z jet tagging in $t\overline{t}$, QCD, and signal simulation for the tZb analysis.

## 4.4 Tagging b jets

The b quark from the VLQ or W′ boson decay is reconstructed as an AK4 jet that is required to pass the DeepJet b tagging algorithm [56, 57] at an operating point defined by a 1% light quark and gluon misidentification probability and approximately 80% b quark efficiency. A simulation-to-data scale factor for the b tagging requirement is used to improve the agreement between data and simulation.

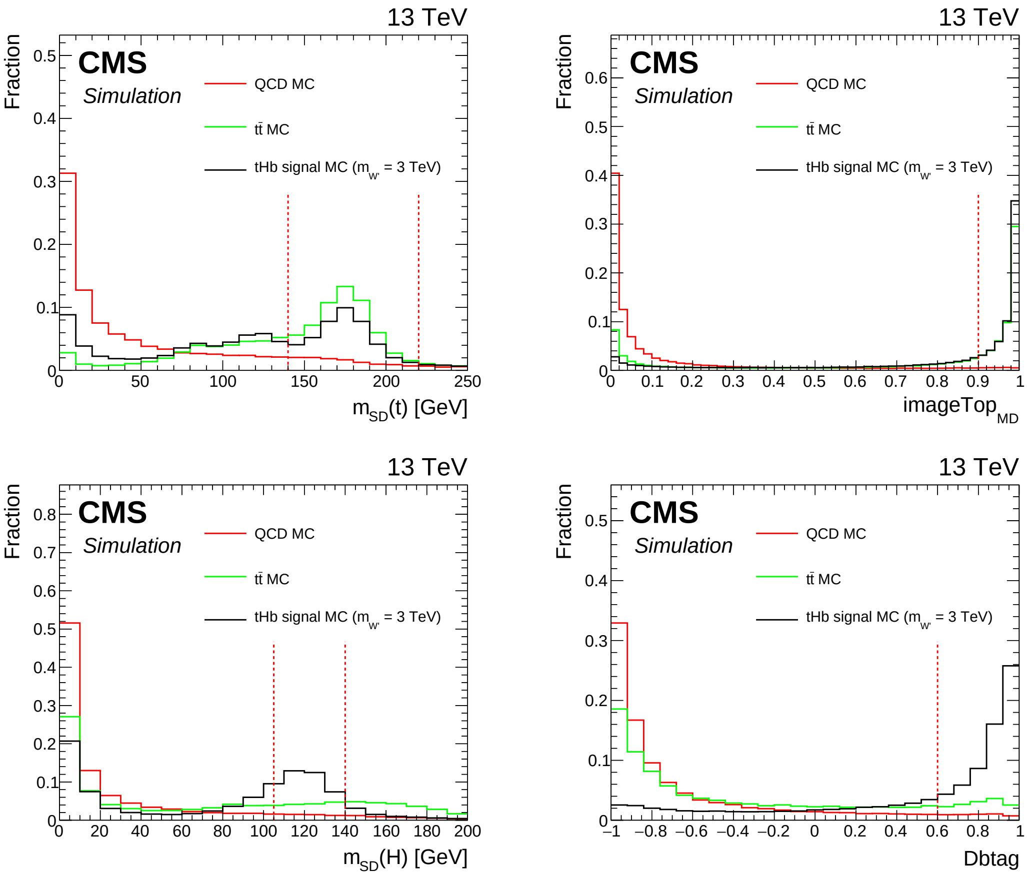


Figure 2: Simulated distributions of the discriminating variables for $t\bar{t}$, QCD, and tHb signal simulated events normalized to unity for the tHb analysis. Discriminant thresholds are shown as vertical dashed lines.

### 4.5 Event selection

We categorize events into selection regions that are either used to search for the W′ boson signal, predict the background in the signal region, or validate this background estimate. Event selection details are presented in Table 1. The signal region is required to contain a tight t jet, a Higgs or Z jet, and a b jet. Regions used for background estimation and validation require combinations of loosened tagger criteria.

The sensitivity of the selections has been studied both in the context of the expected cross section upper limit and of the W′ boson discovery potential. After identifying the t, Higgs or Z, and b jets, the W′ boson candidate mass is taken as the invariant mass of the three jets. Table 2 shows the signal efficiency for all samples considered.

## 5 Background estimation

The primary background is from QCD multijet production and is estimated from data. For this estimate, we use control regions that are selected with identical kinematic criteria to the signal region, but with a reduced acceptance for signal events. Table 1 and Fig. 4 define various selection regions. A transfer function $TF(p_{\mathrm{T}}, \eta)$ is extracted from data by inverting the Higgs or Z jet candidate selection. After this inversion, $TF(p_{\mathrm{T}}, \eta)$ is defined as the ratio of the jet $p_{\mathrm{T}}$ spectrum of the t candidate in the case that it passes (region B) to the case that it fails (region

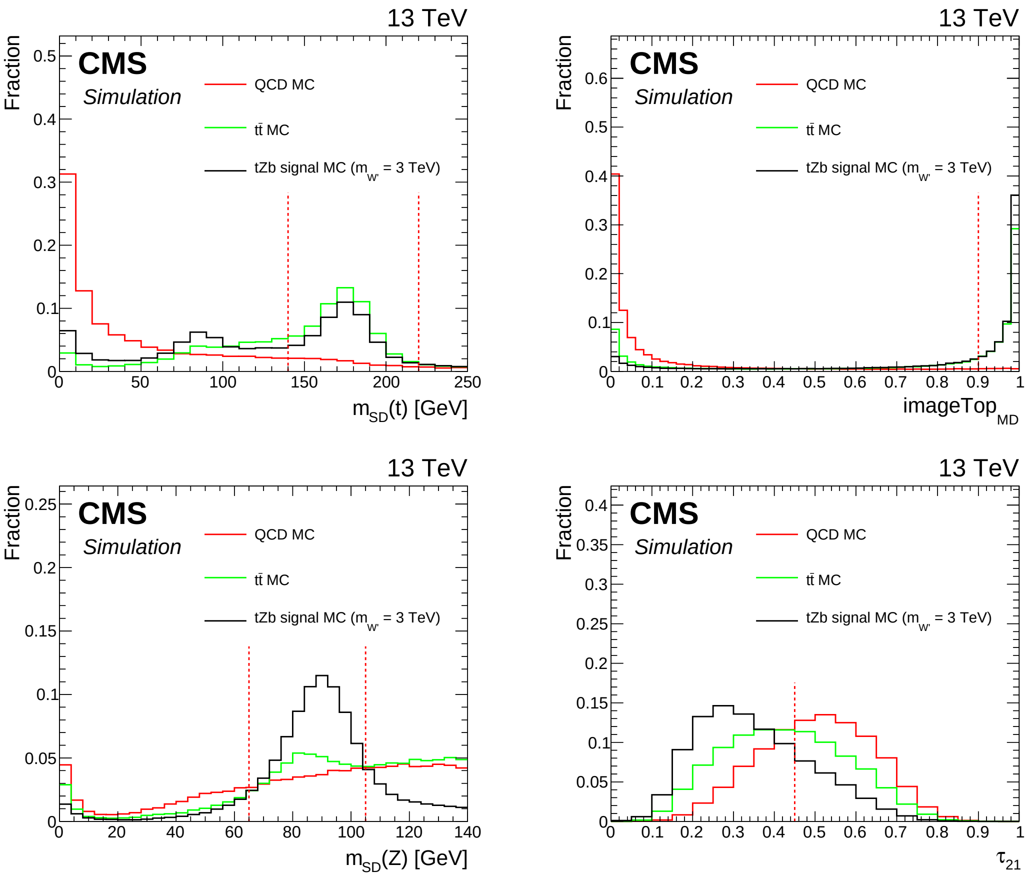


Figure 3: Simulated distributions of the discriminating variables for $t\bar{t}$, QCD, and tZb signal simulated events normalized to unity for the tZb analysis. Discriminant thresholds are shown as vertical dashed lines.

Table 1: Selection regions used for signal identification and background estimation. The AK8 jet discriminant and mass selections are explicitly defined here for the t, Higgs, and Z jet tags.

| Label | Tag | Discriminant | Mass |
|---|---|---|---|
| Tight | H | $0.6 < \text{Dbtag}$ | $105 < m_{SD}(\text{H}) < 140\,\text{GeV}$ |
| | Z | $\tau_{21} < 0.45$ | $65 < m_{SD}(\text{Z}) < 105\,\text{GeV}$ |
| | t | $0.9 < \text{imageTop}_{\text{MD}}$ | $140 < m_{SD}(\text{t}) < 220\,\text{GeV}$ |
| Medium | H | $0 < \text{Dbtag} < 0.6$ | $105 < m_{SD}(\text{H}) < 140\,\text{GeV}$ |
| | Z | $0.45 < \tau_{21} < 0.6$ | $65 < m_{SD}(\text{Z}) < 105\,\text{GeV}$ |
| | t | $0.3 < \text{imageTop}_{\text{MD}} < 0.9$ | $140 < m_{SD}(\text{t}) < 220\,\text{GeV}$ |
| Inverted | H | $-1 < \text{Dbtag} < 0$ | $5 < m_{SD}(\text{H}) < 30\,\text{GeV}$ |
| | Z | $0.6 < \tau_{21} < 1$ | $5 < m_{SD}(\text{Z}) < 30\,\text{GeV}$ |
| | t | $0 < \text{imageTop}_{\text{MD}} < 0.3$ | $30 < m_{SD}(\text{t}) < 65\,\text{GeV}$ |

Table 2: The signal efficiency (in percent) for the three VLQ mass ranges considered. The efficiency is given for the tHb and tZb final states, separated into the low, medium, and high VLQ mass categories.

| $m_{W'}$(GeV) | tHb low | tHb medium | tHb high | tZb low | tZb medium | tZb high |
|---|---|---|---|---|---|---|
| 1500 | 0.29 | 0.34 | 0.28 | 0.43 | 0.51 | 0.39 |
| 2000 | 1.2 | 1.4 | 0.78 | 1.7 | 1.9 | 1.1 |
| 2500 | 2.0 | 1.9 | 1.3 | 2.8 | 2.7 | 1.8 |
| 3000 | 2.3 | 2.2 | 1.8 | 2.7 | 3.0 | 2.5 |
| 3500 | 2.3 | 2.2 | 1.8 | 3.2 | 3.0 | 2.5 |
| 4000 | 2.2 | 2.0 | 1.8 | 3.1 | 2.9 | 2.6 |
| 4500 | 2.0 | 1.9 | 1.7 | 2.8 | 2.7 | 2.5 |
| 5000 | 1.8 | 1.7 | 1.5 | 2.6 | 2.5 | 2.4 |

A) the t tagging algorithm. The $TF(p_T, \eta)$ is extracted in two $\eta$ ranges (central, $|\eta| < 1.2$, and forward, $|\eta| > 1.2$, regions).

The $TF(p_T, \eta)$ distribution is used to predict the background in the signal region. This is accomplished by defining the control region D, which has identical Higgs or Z and b jet candidate selections to the signal region, but with an inverted t jet selection. Events that pass the region D selection use $TF(p_T, \eta)$ as an event weight in a given t-candidate jet ($p_T$, $\eta$) bin. These weighted events are then used to provide the QCD background estimate in the signal region. The primary assumption for this background estimate method is that the t jet substructure selection can be inverted without significantly biasing the Higgs or Z jet substructure selection.

In the $TF(p_T, \eta)$ extraction procedure, the $t\bar{t}$ production component derived from simulation is subtracted from data to ensure that $TF(p_T, \eta)$ refers only to the QCD component. Additionally, the $t\bar{t}$ contamination of the QCD background estimate in the signal region must be subtracted. This is performed by applying the QCD background estimation procedure to simulated $t\bar{t}$ events using the same $TF(p_T, \eta)$ as is used when extracting the QCD estimate from data. The following relations are used to extract the QCD component of the signal region:

$$\begin{aligned} TF(p_T, \eta) &\equiv (B_{data} - B_{t\bar{t}})/(A_{data} - A_{t\bar{t}}), \\ C_{QCD} &\simeq (D_{data} - D_{t\bar{t}})TF(p_T, \eta). \end{aligned} \tag{2}$$

Here, the subscripts "data" and "$t\bar{t}$" indicate that the distributions are obtained from data and $t\bar{t}$ events, respectively.

To test the validity of the background estimate in data, a series of "validation" regions are defined (F, K, and H, as defined in Fig. 4). The transfer function used for predicting the background in the K and F regions ($TF_v(p_T, \eta)$) is estimated from the ratio of regions E to A, whereas the transfer function for the H region is the same as for the signal region. The following relations are used to extract the QCD background component of the validation regions:

$$\begin{aligned} TF_v(p_T, \eta) &\equiv (E_{data} - E_{t\bar{t}})/(A_{data} - A_{t\bar{t}}), \\ H_{QCD} &\simeq (G_{data} - G_{t\bar{t}})TF(p_T, \eta), \\ K_{QCD} &\simeq (D_{data} - D_{t\bar{t}})TF_v(p_T, \eta), \\ F_{QCD} &\simeq (G_{data} - G_{t\bar{t}})TF_v(p_T, \eta). \end{aligned} \tag{3}$$

The background validation tests in the F, K, and H regions are shown in Fig. 5. These validation regions confirm the background estimate agreement with $\chi^2$/ndf values of 1.7, 1.0, and 1.3, respectively for the combined fit which considers systematic uncertainties (see Section 6), where

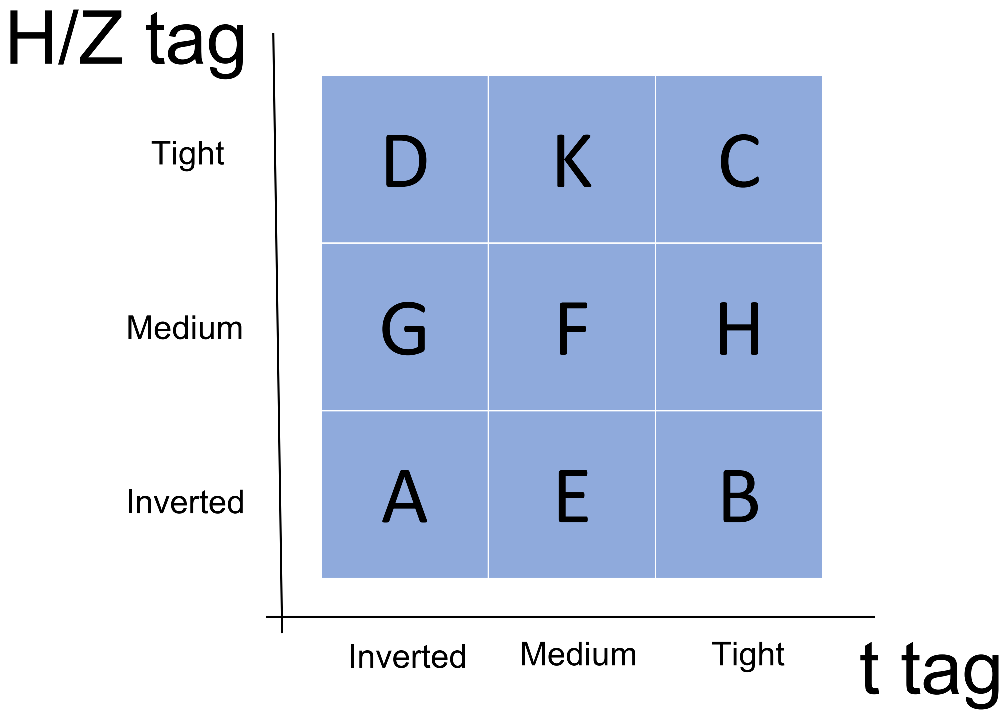


Figure 4: Diagram showing the regions used for background estimation. The signal region is C, the regions A and B are used to determine $TF(p_{\mathrm{T}}, \eta)$, and F, K, and H are validation regions. The $x$ axis indicates the t tag category, and the $y$ axis represents the Higgs or Z boson tag category. The inverted, medium, and tight tag category definitions are given in Table 1.

ndf is the number of degrees of freedom. The $\mathrm{t\bar{t}}$ background component in these validation regions is subtracted using the same procedure that is used when performing the signal region background estimate. The agreement demonstrates that the t jet selection can be inverted without biasing the Higgs jet selection.

The $\mathrm{t\bar{t}}$ background component is estimated using simulated events with an additional event weight that corrects the generator t jet $p_{\mathrm{T}}$ distribution [58]. This correction reduces the disagreement between simulation and data due to a known generator-level mismodeling effect and is validated using a $\mathrm{t\bar{t}}$ control region. The final state for this control region contains two boosted AK8 t jets, so the selections on the AK4 b jet are loosened, and both of the AK8 jets are required to pass the same imageTop$_{\mathrm{MD}}$ selection as the signal region. After this selection, the composition is a nearly pure (approximately 99%) sample of hadronic $\mathrm{t\bar{t}}$ events, demonstrating the discrimination power of the imageTop$_{\mathrm{MD}}$ tagger. The QCD background estimate in this control region is performed in an analogous way to the signal region, and is estimated by weighting events with one inverted and one tight t tag by $TF(p_{\mathrm{T}}, \eta)$. This control region is used to extract an additional correction, which is taken as a simple ratio of QCD-subtracted data to the $\mathrm{t\bar{t}}$ simulation estimate. The uncertainty in this procedure is taken as the uncorrelated combination of the statistical uncertainty in the ratio, 100% of the QCD estimate, and the t tag scale factor uncertainty. This correction is applied to all simulated $\mathrm{t\bar{t}}$ distributions.

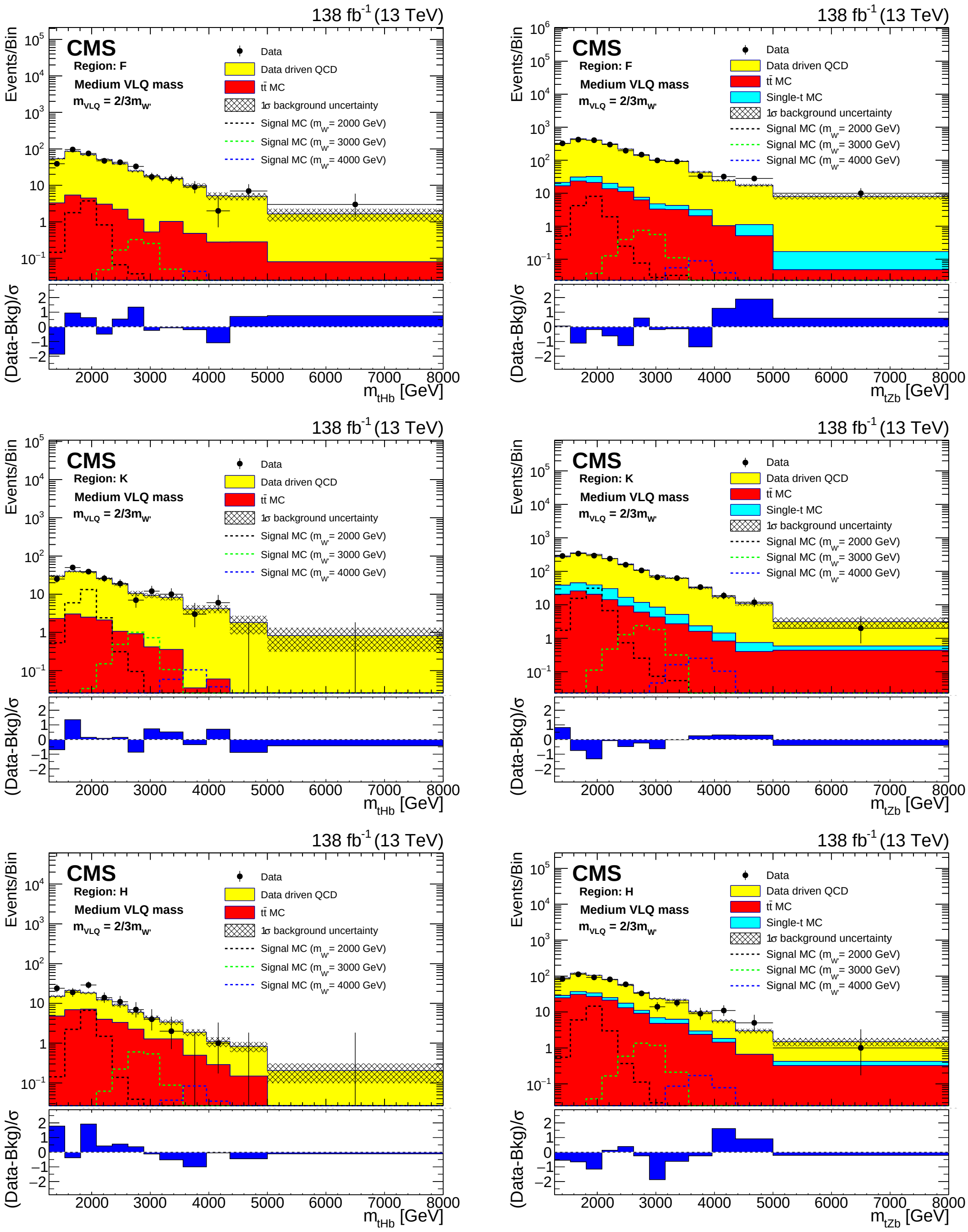


Figure 5: Background closure test for the reconstructed W′ boson invariant mass in region F (upper), K (middle), and H (lower) for the purpose of validation in the tHb (left) and tZb (right) analyses. The data are shown as points with error bars and the estimated backgrounds as solid histograms after a background-only fit, with a hatched band indicating the uncertainty. Expected signals for three different W′ mass hypotheses, assuming the medium VLQ mass, are shown as dotted histograms. The lower panel of each plot shows the difference between the number of events observed in the data and the predicted background, divided by the total uncertainty in the background and the statistical uncertainty in the data added in quadrature.

# 6 Systematic uncertainties

A wide range of systematic uncertainties are considered, and organized into those that impact only the event yields, which are assumed to follow log-normal distributions [59], and those that affect the shape of the three-jet invariant mass distribution as well, which are assumed to follow Gaussian distributions. All of the systematic uncertainties are summarized in Table 3.

Table 3: Sources of systematic uncertainty that affect the final distributions. Sources that affect the normalization only are represented by a range of values for the systematic variation. Sources listing the systematic variation as $\pm 1\sigma(x)$ affect the shape of the three-jet invariant mass distribution as well, which is dependent on $x$.

| Source | Variation | Process |
|---|---|---|
| Integrated luminosity | ±1.2–2.5% | signal, $\mathrm{t\bar{t}}$, single t |
| t jet tagging | $\pm 1\sigma(p_{\mathrm{T}})$ | signal, $\mathrm{t\bar{t}}$, single t |
| JES | $\pm 1\sigma(p_{\mathrm{T}}, \eta)$ | signal, $\mathrm{t\bar{t}}$, single t |
| JER | $\pm 1\sigma(p_{\mathrm{T}}, \eta)$ | signal, $\mathrm{t\bar{t}}$, single t |
| JMS | $\pm 1\sigma(m_{\mathrm{SD}})$ | signal, $\mathrm{t\bar{t}}$, single t |
| JMR | $\pm 1\sigma(m_{\mathrm{SD}})$ | signal, $\mathrm{t\bar{t}}$, single t |
| b tagging | $\pm 1\sigma(p_{\mathrm{T}})$ | signal, $\mathrm{t\bar{t}}$, single t |
| b mistagging | $\pm 1\sigma(p_{\mathrm{T}})$ | signal, $\mathrm{t\bar{t}}$, single t |
| Dbtag tagging | $\pm 1\sigma(p_{\mathrm{T}})$ | signal |
| Dbtag mistagging | $\pm 1\sigma(p_{\mathrm{T}})$ | signal, $\mathrm{t\bar{t}}$, single t |
| Z jet tagging | $\pm 1\sigma(p_{\mathrm{T}})$ | signal |
| Pileup | $\pm 1\sigma$ | signal, $\mathrm{t\bar{t}}$, single t |
| PDF, $\alpha_s$ | $\pm 1\sigma$ | signal, single t |
| $\mu_{\mathrm{R}}$ and $\mu_{\mathrm{F}}$ | $\pm 1\sigma$ | signal, single t |
| ISR/FSR | $\pm 1\sigma$ | single t |
| $\mathrm{t\bar{t}}$ normalization | ±11–15% | $\mathrm{t\bar{t}}$ |
| Trigger | $\pm 1\sigma(H_{\mathrm{T}})$ | signal, $\mathrm{t\bar{t}}$, single t |
| $TF(p_{\mathrm{T}}, \eta)$ | $\pm 1\sigma(p_{\mathrm{T}}, \eta)$ | QCD |
| $\mathrm{t\bar{t}}$ contamination | $\pm 1\sigma(p_{\mathrm{T}}, \eta)$ | QCD |

## 6.1 Normalization uncertainties

The uncertainty in the integrated luminosity is taken as 1.2% for 2016, 2.3% for 2017, and 2.5% for 2018 data [22–24].

The uncertainty in the $\mathrm{t\bar{t}}$ normalization measurement (see Section 5) used to correct the estimate in the signal region accounts for the full uncertainty in the $\mathrm{t\bar{t}}$ yield. It is calculated separately for data collected in 2016, 2017, and 2018, and ranges from 11 to 15%. This is considered a measurement of $\mathrm{t\bar{t}}$ uncertainties that have a large normalization effect, including top $p_{\mathrm{T}}$ reweighting, PDF, and renormalization and factorization ($\mu_{\mathrm{R}}$ and $\mu_{\mathrm{F}}$) scale uncertainties. Uncertainties that have a pronounced shape effect on the invariant mass distribution are included separately when setting limits (see Table 3).

## 6.2 Shape uncertainties

The uncertainty in the JES is taken into account by varying the JES within its uncertainty and observing the effect on the invariant mass distribution. The JES variation impacts the invariant mass distribution shape through a horizontal shift but also causes a normalization difference when the jet falls above or below a kinematic threshold. This uncertainty is propagated to the t jet $m_{SD}(t)$ estimate. The uncertainty in the JER is also taken into account by the uncertainty in the JER correction used for simulated events. This uncertainty is applied to all simulated events and has only a small impact on the final limit.

The uncertainty in the JMS and JMR is measured in a highly enriched sample of $t\bar{t}$ events containing one final-state lepton. In this sample, a fit is performed to the W boson PUPPI $m_{SD}$ distribution, from which the mean and width are extracted. The JMS uncertainty is estimated from the uncertainty in the shift of the W mass peak, and the uncertainty in the JMR is estimated from the uncertainty in the width. These uncertainties are applied to the Higgs or Z boson $m_{SD}$ estimates for all simulated events and have an $\approx$4% effect on the signal estimate.

The uncertainty in the $\tau_{21}$ selection associated with Z jet tagging includes a constant uncertainty and a $p_T$-dependent uncertainty based on an extrapolation to momenta beyond the kinematic region of the scale factor measurement. This uncertainty is applied to all simulated events and has an $\approx$15% effect.

The uncertainty used for the b tagging requirement on the AK4 jet is evaluated by varying the b tagging and b mistagging scale factors within their uncertainties [54]. These uncertainty sources are considered uncorrelated and are applied to all simulated events and result in an $\approx$5% effect on the signal estimate.

The uncertainty in the Dbtag method used for the Higgs jet tagging [54] selection is evaluated by varying the Dbtag scale factor by the uncertainty. The scale factor is parameterized using three regions in $p_T$. Also evaluated is the scale factor for a Higgs boson that is mistagged as a t quark (see Section 4). The uncertainties in both the Higgs jet tagging efficiency and the mistag probability are applied to all simulated events and are treated as uncorrelated during limit setting. The Higgs boson Dbtag uncertainty has an $\approx$10% effect on the signal estimate.

The uncertainty in the imageTop$_{MD}$ selection used for t tagging is $p_T$ dependent and evaluated by varying the imageTop$_{MD}$ scale factor within the uncertainty. This has an $\approx$6% effect on the signal estimate.

The kinematic requirements ensure that the trigger is close to being fully efficient, and all simulated events are corrected using the method outlined in Section 4. The uncertainty in this correction is taken to be half of the trigger inefficiency. This uncertainty is small and is applied to all simulated events.

As mentioned in Section 3, the simulated pileup distribution is reweighted to match data using an effective total inelastic cross section. The uncertainty in this procedure is evaluated by varying the total inelastic cross section by $\pm$4.6% [60]. The resultant uncertainty is applied to all simulated events and has only a small impact on the final limit.

The PDF uncertainty is evaluated using the NNPDF3.1 [42, 43] set for the signal, and the PDF4LHC15 combined set [61, 62] for single t quark events. The uncertainty is calculated using the symmetric Hessian approach.

The uncertainty in the $\mu_R$ and $\mu_F$ scales are evaluated by varying the $\mu_R$ or $\mu_F$ scales up and down by a factor of two. Both independent and simultaneous variations of $\mu_R$ and $\mu_F$ are considered, resulting in six weights after excluding the two most extreme variations. The envelope of these weights is used as the uncertainty. The $\mu_R$, $\mu_F$, and PDF uncertainties are taken as the signal cross section theoretical uncertainty. For the single t quark prediction, these uncertainties are considered when setting limits in addition to the parton shower variation in the initial- and final-state radiation (ISR and FSR). Together, the effect of these theoretical uncertainties amount to $\approx$65%, and dominate the total uncertainty in the single top quark estimate.

The primary uncertainty in the shape of the QCD background estimate is taken from the statistical uncertainty in the $TF(p_{\mathrm{T}}, \eta)$. This is propagated to the invariant mass spectrum by evaluating the $TF(p_{\mathrm{T}}, \eta)$ weight at $\pm 1\sigma$ in a given ($p_{\mathrm{T}}$, $\eta$) bin. The uncertainty from each $TF(p_{\mathrm{T}}, \eta)$ bin is added in quadrature to form the full uncertainty, and the effect on the QCD background estimate is $\approx$20%. The effect of the uncertainty in the $\mathrm{t\bar{t}}$ subtraction procedure is evaluated by varying the amount subtracted within the range of the $\mathrm{t\bar{t}}$ normalization uncertainty. The effect on the QCD background estimate is $\approx$6%.

The statistical uncertainty of the simulation is taken into account during limit setting by using the "Barlow–Beeston lite" method [63].

## 7 Results

The final distributions of reconstructed $m_{\mathrm{W'}}$ for the signal region are shown in Fig. 6. Since no deviation from the SM expectation is observed, we place limits on the W′ boson production cross section.

The 95% confidence level (CL) upper limits are obtained using the $\mathrm{CL_s}$ approach [64, 65] in the asymptotic approximation [66] on the product of the W′ boson production cross section for the $s_{\mathrm{L}} = 0.5$ and $\cot(\theta_2) = 3$ hypothesis [13], and for the benchmark branching fraction. A binned maximum likelihood fit to the data is used to calculate upper limits, in a process where all systematic uncertainties are included as nuisance parameters. For the signal template, the sum of the invariant mass distributions from the tB and bT decay channels are used. Upper limits on the product of cross section and branching fraction are shown in Fig. 7. A W′ boson with a mass below 3.1 TeV is excluded at 95% CL for the medium VLQ mass case. The low and high VLQ mass benchmarks have a lower $\mathrm{W'} \to$ VLQ branching fraction, and the sensitivity is not sufficient to set mass exclusion limits.

Variations of the primary assumptions of this benchmark are shown in Fig. 8 for the purpose of generalizing the results. Specifically, the fraction of bT and tB from the W′ decay is varied between 0 and 1, and the VLQ branching fractions to qH and qZ are varied within the same bounds. Therefore, for both generalized coupling figures, the mass exclusion of the benchmark model corresponds to the point with coordinates $[0.5, 0.5]$, indicated with an asterisk.

## 8 Summary

A search has been presented for a heavy W′ boson decaying to a B or T vector-like quark and a t or b quark, respectively. The data correspond to an integrated luminosity of 138 $\mathrm{fb^{-1}}$ collected between 2016 and 2018 with the CMS detector at the LHC in proton-proton collisions at $\sqrt{s} = 13$ TeV. The signature considered for both decay modes is a t quark and a Higgs or Z boson, each decaying hadronically, and a b quark jet. Boosted heavy-resonance identification techniques are used to select the events containing three energetic jets and to suppress standard model backgrounds. No significant deviation from the standard model background prediction is observed. Upper limits are placed on the product of the W′ boson cross section and the final state branching fraction as a function of the $m_{\mathrm{W'}}$. A W′ boson with a mass below 3.1 TeV is excluded at 95% confidence level, given the benchmark model assumption of democratic branching fractions. In addition, limits are set based on generalizations of these assumptions. These are the most sensitive limits to date for this final state.

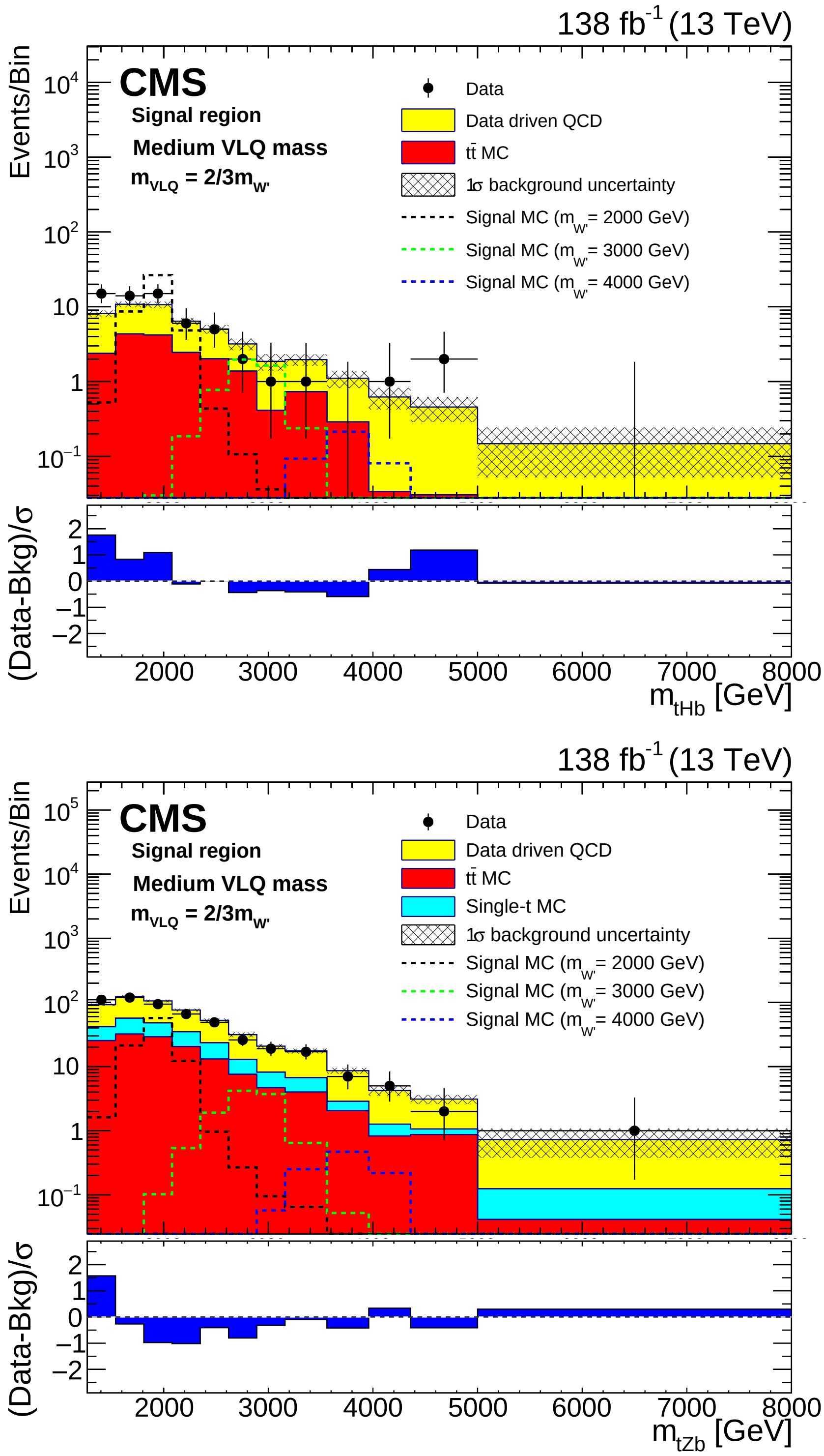


Figure 6: Reconstructed $m_{W'}$ distributions ($m_{tHb}$ (upper), and $m_{tZb}$ (lower)) in the signal region with estimated backgrounds and signal for several W′ boson mass hypotheses at the medium VLQ mass, after a background-only fit. The combined statistical and systematic uncertainty in the total estimated background is indicated by the hatched region. The lower panels show the difference between the number of events observed in the data and the predicted background, divided by the total uncertainty in the background and the statistical uncertainty in the data added in quadrature.

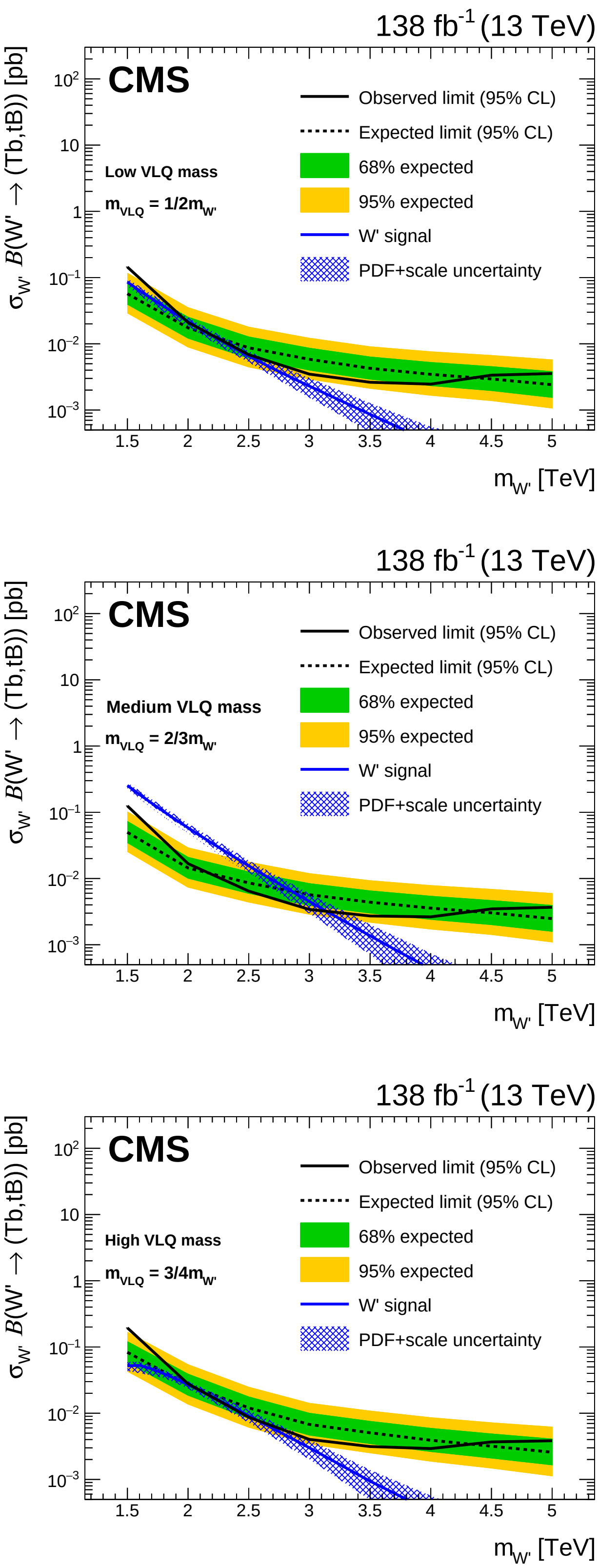


Figure 7: The W′ boson 95% CL limits on the product of cross section and branching fraction. The expected (dashed) and observed (solid) limits, as well as the W′ boson theoretical cross section, with its PDF and scale normalization uncertainties, are shown. The green (inner) and yellow (outer) bands indicate the 68% and 95% confidence intervals of the expected limit. The limits are given for the low (upper), medium (center), and high (lower) VLQ mass ranges.

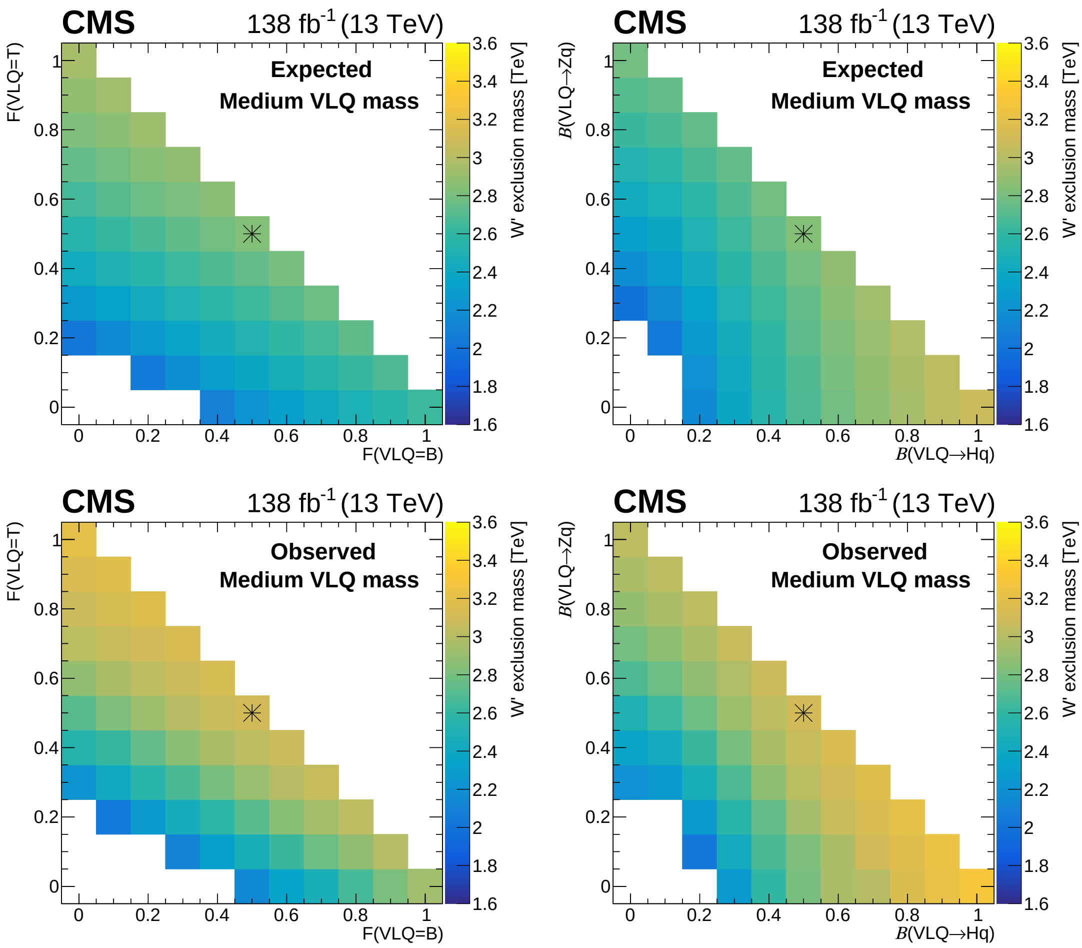


Figure 8: Expected (upper) and observed (lower) 95% CL limits for generalized hypotheses varying the fraction of tB (F(VLQ=B)) and bT (F(VLQ=T)) from the W′ decay (left), and the VLQ branching fraction to qH and qZ (right). The asterisk marker signifies the branching fractions for the benchmark model

## Acknowledgments

We congratulate our colleagues in the CERN accelerator departments for the excellent performance of the LHC and thank the technical and administrative staffs at CERN and at other CMS institutes for their contributions to the success of the CMS effort. In addition, we gratefully acknowledge the computing centers and personnel of the Worldwide LHC Computing Grid and other centers for delivering so effectively the computing infrastructure essential to our analyses. Finally, we acknowledge the enduring support for the construction and operation of the LHC, the CMS detector, and the supporting computing infrastructure provided by the following funding agencies: BMBWF and FWF (Austria); FNRS and FWO (Belgium); CNPq, CAPES, FAPERJ, FAPERGS, and FAPESP (Brazil); MES and BNSF (Bulgaria); CERN; CAS, MoST, and NSFC (China); MINCIENCIAS (Colombia); MSES and CSF (Croatia); RIF (Cyprus); SENESCYT (Ecuador); MoER, ERC PUT and ERDF (Estonia); Academy of Finland, MEC, and HIP (Finland); CEA and CNRS/IN2P3 (France); BMBF, DFG, and HGF (Germany); GSRI (Greece); NKFIA (Hungary); DAE and DST (India); IPM (Iran); SFI (Ireland); INFN (Italy); MSIP and NRF (Republic of Korea); MES (Latvia); LAS (Lithuania); MOE and UM (Malaysia); BUAP, CINVESTAV, CONACYT, LNS, SEP, and UASLP-FAI (Mexico); MOS (Montenegro); MBIE (New Zealand); PAEC (Pakistan); MSHE and NSC (Poland); FCT (Portugal); JINR (Dubna); MON, RosAtom, RAS, RFBR, and NRC KI (Russia); MESTD (Serbia); MCIN/AEI and PCTI (Spain); MOSTR (Sri Lanka); Swiss Funding Agencies (Switzerland); MST (Taipei); ThEPCenter, IPST, STAR, and NSTDA (Thailand); TUBITAK and TAEK (Turkey); NASU (Ukraine); STFC (United Kingdom); DOE and NSF (USA).

Individuals have received support from the Marie-Curie program and the European Research Council and Horizon 2020 Grant, contract Nos. 675440, 724704, 752730, 758316, 765710, 824093, 884104, and COST Action CA16108 (European Union); the Leventis Foundation; the Alfred P. Sloan Foundation; the Alexander von Humboldt Foundation; the Belgian Federal Science Policy Office; the Fonds pour la Formation à la Recherche dans l'Industrie et dans l'Agriculture (FRIA-Belgium); the Agentschap voor Innovatie door Wetenschap en Technologie (IWT-Belgium); the F.R.S.-FNRS and FWO (Belgium) under the "Excellence of Science – EOS" – be.h project n. 30820817; the Beijing Municipal Science & Technology Commission, No. Z191100007219010; the Ministry of Education, Youth and Sports (MEYS) of the Czech Republic; the Deutsche Forschungsgemeinschaft (DFG), under Germany's Excellence Strategy – EXC 2121 "Quantum Universe" – 390833306, and under project number 400140256 - GRK2497; the Lendület ("Momentum") Program and the János Bolyai Research Scholarship of the Hungarian Academy of Sciences, the New National Excellence Program ÚNKP, the NKFIA research grants 123842, 123959, 124845, 124850, 125105, 128713, 128786, and 129058 (Hungary); the Council of Science and Industrial Research, India; the Latvian Council of Science; the Ministry of Science and Higher Education and the National Science Center, contracts Opus 2014/15/B/ST2/03998 and 2015/19/B/ST2/02861 (Poland); the Fundação para a Ciência e a Tecnologia, grant CEECIND/01334/2018 (Portugal); the National Priorities Research Program by Qatar National Research Fund; the Ministry of Science and Higher Education, projects no. 0723-2020-0041 and no. FSWW-2020-0008 (Russia); MCIN/AEI/10.13039/501100011033, ERDF "a way of making Europe", and the Programa Estatal de Fomento de la Investigación Científica y Técnica de Excelencia María de Maeztu, grant MDM-2017-0765 and Programa Severo Ochoa del Principado de Asturias (Spain); the Stavros Niarchos Foundation (Greece); the Rachadapisek Sompot Fund for Postdoctoral Fellowship, Chulalongkorn University and the Chulalongkorn Academic into Its 2nd Century Project Advancement Project (Thailand); the Kavli Foundation; the Nvidia Corporation; the SuperMicro Corporation; the Welch Foundation, contract C-1845; and the Weston Havens Foundation (USA).

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